# The Grüneisen Parameter and Higher Order Thermo Elastic Properties of Fe, Co at Higher Pressures

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#### **ABSTRACT**

Expressions have been obtained for the higher order Thermo elastic properties of the Grüneisen parameter in the limit of infinite pressure The usefulness of these expressions has been demonstrated by considering a relationship between reciprocal Grüneisen parameter and the pressure bulk modulus ratio at finite pressure. The results have been obtained using the shanker reciprocal gamma relationship. The results have also been determined using the Stacey-Davis formulations for the third order Grüneisen parameters. The calculations have been performed in case of transition metals Fe & Co for the wide range of compressions. It has been found that both the models yield close agreement with each other. The obtained results in the present paper satisfy the thermodynamic constraints in the limit of infinite pressure.

KEYWORDS: Grüneisen parameters, Thermoelastic properties, infinite pressure behavior, transition metal

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# INTRODUCTION

The Grüneisen parameter  $\gamma$  is related to Thermoelastic properties of materials [1-3]. The pressure derivatives of  $\gamma$ play a central role in predicting the thermal behavior, equation of state, and melting at high pressures [4-6]. The Grüneisen parameter  $\gamma$  is an important physical quantity directly related to thermal and elastic properties of

materials [7–9] as follows
$$\gamma = \frac{\alpha K_T V}{c_V} = \frac{\alpha K_S V}{c_P}$$
(1)

Where  $\alpha$  is the thermal expansivity  $K_T$  and  $K_S$  are isothermal and adiabatic bulk moduli,  $C_V$  and Cp are specific heats at constant volume and constant pressure, respectively. For a given material y decreases with increasing pressure P or decreasing volume V.

There have been various attempts [7-11] to formulate expressions for  $\gamma(V)$  as well as  $\gamma(P)$ . The most critical examination of these formulations can be made by applying the thermodynamic constraints for the higherorder Grüneisen parameters [7,9, 15-17] at extreme compression in the limit of infinite pressure. In the presented study, we have been determined the higher order Grüneisen parameter in case of transition metals Fe, Co for the wide range of compressions.

# Formulation for $\gamma(V)$ :

According to the Stacey- Davis model we can write 
$$\lambda = \lambda_{\infty} + (\lambda_0 - \lambda_{\infty}) \frac{q}{q_0}$$
 (2)

On integrating Eq. (2), we get

$$q = q_0 \left[ 1 + \frac{\lambda_0}{\lambda_\infty} \left( \frac{V_0}{V} \right)^{\lambda_\infty} - \frac{\lambda_0}{\lambda_\infty} \right]^{-1}$$
 (3)

And further integration gives

$$\gamma = \gamma_0 \left[ \frac{\lambda_0}{\lambda_\infty} - \left( \frac{\lambda_0}{\lambda_\infty} - 1 \right) \left( \frac{\nu_0}{\nu} \right)^{\lambda_\infty} \right]^{-\frac{q_0}{(\lambda_0 - \lambda_\infty)}} \tag{4}$$

It has been found [5,8] that if we assume

$$(\lambda_0 - \lambda_\infty) = q_0 \tag{5}$$

Eq. (13) On differentiation yields

$$\xi = \left(\frac{\lambda_0 - \lambda_\infty}{q_0}\right) q \tag{6}$$

In the present study we use Eq. (2), (3), (4) and (6) to obtain  $\gamma$ , q,  $\lambda$  and  $\xi$  for transition metals.

# Formulation for $\gamma(P)$

The generalized free -volume formula for γ as a function of pressure P, bulk modulus and its pressure derivative K' = dK / dP has been widely used in literature [18–21]. This formula is written as follows [8-10]

$$\gamma = \frac{\frac{K}{2} - \frac{1}{6} \cdot \frac{f}{3} \left(1 - \frac{P}{3K}\right)}{\left(1 - 2f \cdot \frac{P}{3K}\right)} \tag{7}$$

where f is the free –volume parameter. Different values of fsuch as 0, 1, 2 and 2.35 have been taken by Slater [18],

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Vashchenko and Zubarev [20], and by Barton and Stacey [21]. The main difficulty associated with Eq. (7) is regarding the uncertainties in the value of f arising from various approximations involved in the derivation of the formula [9]. It is even more uncertain to use Eq. (7) for evaluating derivatives of  $\gamma$  because of difficulties in determining df / dP and d2 f / dP2 . Such calculations are tedious at finite pressures. However, it should be indicating that Eq. (7) in the limit of infinite pressure gives expressions [9, 15] which are free from the free -volume parameter f. An alternative formulation for \_(P) has been presented by Shankar et al. [17]. They proposed the following relationship

$$\frac{1}{\gamma} = \frac{1}{\gamma_0} + K_{\infty} \left( \frac{1}{\gamma_{\infty}} - \frac{1}{\gamma_0} \right) \frac{P}{K}$$
 (8)

Where  $K_{\infty}$  is the minimum value of K in the limit of infinite pressure. Eq. (8) satisfies the following identity [20, 3]

$$\left(\frac{P}{K}\right)_{\infty} = \frac{1}{K_{\infty}} \tag{9}$$

Eq. (8) yields the following expressions for the higher order derivatives of the Grüneisen parameter [16]

$$\frac{q}{\gamma} = K_{\infty} \left( \frac{1}{\gamma_{\infty}} - \frac{1}{\gamma_0} \right) \left( 1 - K \frac{P}{K} \right) \tag{10}$$

$$\lambda = \frac{KK'}{(1 - K'P/K)} \frac{P}{K} + K' + q \tag{11}$$

$$\xi = -\frac{2KK'}{\lambda (1 - K'P/K)} - \frac{(K'KP/K)^2}{\lambda (1 - K'P/K)^2} - \frac{K^2K'(\frac{P}{K})}{\lambda (1 - K'P/K)} + q \tag{12}$$

It should be mentioned that the expressions for  $\gamma$ , q,  $\lambda$  and  $\xi$  given above satisfy the boundary conditions in the limit of infinite pressure viz.  $\gamma_{\infty}$  and  $\lambda_{\infty}$  both remain positive finite whereas  $q \infty$  and  $x \infty$  both tend to zero. Equations (8), (10), (11) and (12) at infinite pressure become identical to the corresponding expressions derived from the basic principles of calculus which have the status of identities [13, 14, 16].

### INFINITE PRESSURE BEHAVIOUR

P is a function of V such that increases continuously with the decreasing V. At extreme compression  $(V \rightarrow 0)$ , P becomes infinitely large, then we can write

$$\left(\frac{dlnP}{dlnV}\right)_{V\to 0} = negative finite$$
 (13)

Since 
$$\frac{dlnP}{dlnV} = -\frac{K}{P}$$
 (14)

Where K is bulk modulus

$$K = -V \frac{dP}{dV} \tag{15}$$

Now taking  $y = \frac{K}{P}$  and x = V

in expression 
$$\left(\frac{dlny}{dlnx}\right)_{x\to 0} = 0$$

Where y is a function of x such that y remains positive finite at  $x \rightarrow 0$ 

if y becomes zero at 
$$x \to 0$$

$$\left[\frac{dln(K/P)}{dlnV}\right]_{V \to 0} = 0$$
(17)

Because K/P is positive finite (Eq. 16), on solving Eq. (17),

$$(1 - K' P/K)_{\infty} = 0 (18)$$

Eq. (18) yields 
$$K_{\infty}^{\cdot} = \left(\frac{K}{P}\right)_{\infty} \tag{19}$$

Equation (19) has the status of an identity [7,5, 20]. In view of Eq. (16), it is revealed from Eq. (19) that  $K_{\infty}$ , the pressure derivative of bulk modulus, dK/dP, at infinite pressure remains positive finite.

#### Result and discussion

The quantities calculated as a function of pressure can be converted to those as a function of volume, and vice versa with the help of an equation of state theory. We use the Stacey K- primed equation of state [9] which is based on

Eq.(9). And can be expressed as 
$$\frac{1}{K} = \frac{1}{K_0} + \left(1 - \frac{K_\infty}{K_0}\right) \frac{P}{K}$$
 (20)

On integration, Eq. (20) gives
$$\frac{K}{K_0} = \left(1 - K_{\infty}^{\cdot \cdot} \frac{P}{K}\right)^{-K_0^{\cdot}/K_{\infty}^{\cdot}}$$
(21)

(12) Further integration of Eq. (21) gives 
$$\ln\left(\frac{V}{V_0}\right) = \left(\frac{K_0'}{K_\infty'} - 1\right) \left(\frac{P}{K}\right) + \frac{K_0'}{K_\infty'^2} \ln\left(1 - K_\infty' \frac{P}{K}\right)$$
(22)

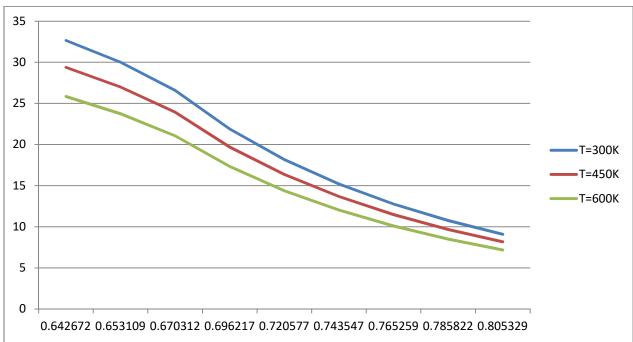
The Grüneisen parameter values  $\gamma$  and its higher order derivatives  $q \lambda$  and  $\xi$  calculated from the Shanker formulation [16] using Eq. (8), (10), (11) and (12). Values of  $\gamma \neq \lambda$  and  $\xi$  get decrease with the increment in value of pressure. In the Stacey-Davis formulation they are expressed as a function of  $V/V_0$ . At extreme compression, V tends to zero,  $q_{\infty}$  and  $\xi_{\infty}$  both tends to zero, but  $\gamma_{\infty}$  and  $\lambda_{\infty}$  both remains positive finite.

#### Conclusion

The Grüneisen parameter at extreme compression, the pressure for the earth inner core and outer core boundary, the melting temperature for Fe and Co is very close. The theory of materials at  $(\infty)$  infinite pressure originally due to Stacey [5,7] and further developed by Shanker et al [12,13] has been extended in the present study. Values of  $\gamma q \lambda$  and  $\xi$  at finite pressures for transition metals (Fe, Co) and Calculated from the two formulations compare with each other. Results are quite similar for different transition metals (Fe, Co). Results can be produce using the input parameter given in table first. It should be emphasized that the pressure derivative of bulk modulus play the central role in determining Thermoelastic properties and equation of state of materials [1,3,21-23]. We thus obtain the following generalized expression for the higher pressure derivative of bulk modulus at infinite pressure.

Table (1) Values of input parameters for NaCl

| Element | Z  | γ <sub>0</sub> | $K_0$ | $K_0$ | $K_{\infty}$ | $\gamma_{\infty}$ | $q_0$ | $\lambda_0$ | $\xi_0$ |
|---------|----|----------------|-------|-------|--------------|-------------------|-------|-------------|---------|
| Fe(bcc) | 26 | 1.83           | 187.6 | 5.8   | 3.48         | 1.33              | 1.13  | 6.13        | 4.39    |
| Co(fcc) | 27 | 1.77           | 216.5 | 4.7   | 2.82         | 1.24              | 1.21  | 5.91        | 4.20    |



Plots between pressure P and compression V/V<sub>0</sub> for Fe & Co at different temperatures

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